

Reduced form of Yang-Mills equation of $SU(3)$ on $R^{4,0}$

Shi He and Wu Yuchun¹⁾

Abstract. The Yang-Mills equation of $SU(3)$ on $R^{4,0}$ is studied in this note. An explicit expression of a differential linear transformation is given. The Yang-Mills equation of $SU(3)$ on Euclidean four-dimensional flat space becomes a reduced form via this differential linear transformation. And then, some solutions of the Yang-Mills equation could to be obtained very easily.

1. Notations

Let $R^{4,0}$ be the Euclidean four-dimensional flat space and $x = (x^1, x^2, x^3, x^4)$ be the coordinates. The matrix $ds^2 = \sum g^{ij} dx^i dx^j$ of $R^{4,0}$ is given by $g^{ii} = 1, i = 1, 2, 3, 4, g^{ij} = 0, i \neq j$. Let $gl(3)$ be the Lie algebra of group $SU(3)$. The anti-Hermitian representations of a basis of $gl(3)$ are taken as follows

$$\begin{aligned} X_1 &= \begin{pmatrix} \mathbf{i} & 0 & 0 \\ 0 & -\mathbf{i} & 0 \\ 0 & 0 & 0 \end{pmatrix}, & X_2 &= \begin{pmatrix} 0 & \mathbf{i} & 0 \\ \mathbf{i} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, & X_3 &= \begin{pmatrix} 0 & 0 & \mathbf{i} \\ 0 & 0 & 0 \\ \mathbf{i} & 0 & 0 \end{pmatrix}, \\ X_4 &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & \mathbf{i} \\ 0 & \mathbf{i} & 0 \end{pmatrix}, & X_5 &= \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, & X_6 &= \begin{pmatrix} 0 & 0 & -1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}, \\ X_7 &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix}, & X_8 &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & \mathbf{i} & 0 \\ 0 & 0 & -\mathbf{i} \end{pmatrix}, \end{aligned}$$

in which $\mathbf{i} = \sqrt{-1}$. There are several choices of the base X_8 . For convenient to compute by using computer, we set X_8 as above.

There are 32 unknown functions $A^{ij}(x), i = 1, 2, \dots, 8, j = 1, 2, 3, 4$. Their connections are given by

$$G_j = \sum_{i=1}^8 A^{ij}(x) X_i, \quad j = 1, 2, 3, 4.$$

That is

$$G_j = \begin{pmatrix} \mathbf{i} A^{1j} & \mathbf{i} A^{2j} - A^{5j} & \mathbf{i} A^{3j} - A^{6j} \\ \mathbf{i} A^{2j} + A^{5j} & -\mathbf{i} A^{1j} + \mathbf{i} A^{8j} & \mathbf{i} A^{4j} - A^{7j} \\ \mathbf{i} A^{3j} + A^{6j} & \mathbf{i} A^{4j} + A^{7j} & -\mathbf{i} A^{8j} \end{pmatrix},$$

¹⁾ Institute of Systems Science China Academy of Science Beijing 100080, P.R.China

$$j = 1, 2, 3, 4.$$

The curvatures are given by

$$F_{jk} = \partial_k G_j - \partial_j G_k + G_j G_k - G_k G_j, \quad j, k = 1, 2, 3, 4,$$

in which the notations $\partial_j = \partial/\partial x^j$, $j = 1, 2, 3, 4$.

The Yang-Mills equation of $SU(3)$ on $R^{4,0}$ is

$$\sum_{k=1}^4 g^{kk} (\partial_k F_{jk} + G_k F_{jk} - F_{jk} G_k) = 0, \quad j = 1, 2, 3, 4. \quad (1.1)$$

Note that the Yang-Mills equation in formula (1.1) is a 3×3 matrix:

$$\begin{pmatrix} Ym_{j1} & Ym_{j2} - Ym_{j3} & Ym_{j4} - Ym_{j5} \\ Ym_{j2} + Ym_{j3} & Ym_{j1} + Ym_{j9} & Ym_{j7} - Ym_{j8} \\ Ym_{j4} + Ym_{j5} & Ym_{j7} + Ym_{j8} & Ym_{j9} \end{pmatrix}.$$

2. Differential linear transformation

Using the differential linear transformation, we obtain the reduced form of Yang-Mills equations of $SU(2)$ on Euclidean space $R^{4,0}$ (see [6]), Mincovski space $R^{3,1}$ (see [4]) and Artin space $R^{2,2}$ (see [5]). The goal of this note is to give the reduced form of Yang-Mills equation of $SU(3)$ on Euclidean flat space $R^{4,0}$. Similarly to [6], we have

Definition a transformation $S : u(x) \mapsto v(x)$, $v = C(u)$ between functions $u(x)$ and $v(x)$ is called a **differential linear transformation** if $C(u)$ consists of linear differential operators only.

In [7], Maxwell equation is studied by using differential linear transformation. Now we set

$$A^{ij}(x) = \sum_{k=1}^4 s_{jk} \partial_k u^i(x), \quad i = 1, 2, \dots, 8, \quad j = 1, 2, 3, 4, \quad (1.2)$$

in which $u^i(x)$ are eight new functions and s_{jk} , $j, k = 1, 2, 3, 4$ are undetermined constants. We denote them by a matrix

$$S = \begin{pmatrix} s_{11} & s_{12} & s_{13} & s_{14} \\ s_{21} & s_{22} & s_{23} & s_{24} \\ s_{31} & s_{32} & s_{33} & s_{34} \\ s_{41} & s_{42} & s_{43} & s_{44} \end{pmatrix}.$$

3. Reduced form of Yang-Mills equation

For $R^{4,0}$, we substitute $A^{ij}(x)$, $i = 1, 2, \dots, 8$, $j = 1, 2, 3, 4$ which are given by formula (1.2) in Yang-Mills equation (1.1) and obtain certain product terms such as

$$\partial_j u^s(x) \partial_k u^t(x) \partial_l u^w(x), \quad j, k, l = 1, 2, 3, 4, \quad s, t, w = 1, 2, \dots, 6$$

The coefficients of these terms are polynomials of undetermined constants s_{jk} , $j, k = 1, 2, 3, 4$. We set these polynomials equal to zeros and then obtain a system of polynomial equations. In paper [6] by using Wu elimination [1,2], we get the matrix as following

$$S_{40} = s \begin{pmatrix} 1 & \mathbf{i} & -1 & \mathbf{i} \\ -\mathbf{i} & 1 & -\mathbf{i} & -1 \\ 1 & \mathbf{i} & 1 & -\mathbf{i} \\ -\mathbf{i} & 1 & \mathbf{i} & 1 \end{pmatrix} \tag{1.3}$$

in which s is a parameter. The main conclusion of this note is

Theorem By the differential linear transformation given by formula (1.2) and (1.3), the Yang-Mills equation (1.1) turns out to be a reduced form

$$YM_{ij} = \sum_{k=1}^8 L_{ijk} D_{40}(u^k), \quad i = 1, 2, 3, 4, \quad j = 1, 2, \dots, 8 \tag{1.4}$$

in which

$$D_{40} = \partial_1^2 + \partial_2^2 + \partial_3^2 + \partial_4^2$$

is the Laplace operator and L_{ijk} $i = 1, 2, 3, 4, \quad j = 1, 2, \dots, 8, \quad k = 1, 2, \dots, 8$ are certain linear differential operators. For example, for $i = 1$ we have

$$\begin{aligned} YM_{11} = & s(-\mathbf{i}\partial_2 + \partial_3 - \mathbf{i}\partial_4)D_{40}(u^1) \\ & - 2s^2 \mathbf{i} L_{40}(u^5)D_{40}(u^2) - 2s^2 \mathbf{i} L_{40}(u^6)D_{40}(u^3) \\ & + 2s^2 \mathbf{i} L_{40}(u^2)D_{40}(u^5) + 2s^2 \mathbf{i} L_{40}(u^3)D_{40}(u^6) \end{aligned}$$

$$\begin{aligned} YM_{12} = & s(-\mathbf{i}\partial_2 + \partial_3 - \mathbf{i}\partial_4)D_{40}(u^2) \\ & - 2s^2 L_{40}(u^5)D_{40}(u^1) + s^2 L_{40}(u^7)D_{40}(u^3) \\ & + s^2 L_{40}(u^6)D_{40}(u^4) - s^2 L_{40}(u^4)D_{40}(u^6) \\ & - s^2 L_{40}(u^3)D_{40}(u^7) + s^2 L_{40}(u^5)D_{40}(u^8) \\ & + s^2(2L_{40}(u^1) - L_{40}(u^8))D_{40}(u^5) \end{aligned}$$

$$\begin{aligned} YM_{13} = & s(-\mathbf{i}\partial_2 + \partial_3 - \mathbf{i}\partial_4)D_{40}(u^3) \\ & - s^2 L_{40}(u^6)D_{40}(u^1) - s^2 L_{40}(u^7)D_{40}(u^2) \\ & + s^2 L_{40}(u^5)D_{40}(u^4) - s^2 L_{40}(u^4)D_{40}(u^5) \\ & + s^2 L_{40}(u^2)D_{40}(u^7) - s^2 L_{40}(u^6)D_{40}(u^8) \\ & + s^2(L_{40}(u^1) + L_{40}(u^8))D_{40}(u^6) \end{aligned}$$

$$\begin{aligned} YM_{14} = & s(-\mathbf{i}\partial_2 + \partial_3 - \mathbf{i}\partial_4)D_{40}(u^4) \\ & + s^2 L_{40}(u^7)D_{40}(u^1) - s^2 L_{40}(u^6)D_{40}(u^2) \\ & - s^2 L_{40}(u^5)D_{40}(u^3) + s^2 L_{40}(u^3)D_{40}(u^5) \\ & + s^2 L_{40}(u^2)D_{40}(u^6) - 2s^2 L_{40}(u^7)D_{40}(u^8) \\ & + s^2(2L_{40}(u^8) - L_{40}(u^1))D_{40}(u^7) \end{aligned}$$

$$\begin{aligned} YM_{15} = & s(-\mathbf{i}\partial_2 + \partial_3 - \mathbf{i}\partial_4)D_{40}(u^5) \\ & + 2s^2 L_{40}(u^2)D_{40}(u^1) + s^2 L_{40}(u^4)D_{40}(u^3) \\ & - s^2 L_{40}(u^3)D_{40}(u^4) + s^2 L_{40}(u^7)D_{40}(u^6) \\ & - s^2 L_{40}(u^6)D_{40}(u^7) - s^2 L_{40}(u^2)D_{40}(u^8) \\ & + s^2(2L_{40}(u^1) - L_{40}(u^8))D_{40}(u^2) \end{aligned}$$

$$\begin{aligned}
YM_{16} &= s(-\mathbf{i}\partial_2 + \partial_3 - \mathbf{i}\partial_4)D_{40}(u^6) \\
&\quad + s^2L_{40}(u^3)D_{40}(u^1) + s^2L_{40}(u^4)D_{40}(u^2) \\
&\quad - s^2L_{40}(u^2)D_{40}(u^4) - s^2L_{40}(u^7)D_{40}(u^5) \\
&\quad + s^2L_{40}(u^5)D_{40}(u^7) + s^2L_{40}(u^3)D_{40}(u^8) \\
&\quad - s^2(L_{40}(u^1) + L_{40}(u^8))D_{40}(u^3) \\
YM_{17} &= s(-\mathbf{i}\partial_2 + \partial_3 - \mathbf{i}\partial_4)D_{40}(u^7) \\
&\quad + s^2L_{40}(u^4)D_{40}(u^1) - s^2L_{40}(u^3)D_{40}(u^2) \\
&\quad + s^2L_{40}(u^2)D_{40}(u^3) - s^2L_{40}(u^6)D_{40}(u^5) \\
&\quad + s^2L_{40}(u^5)D_{40}(u^6) - 2s^2L_{40}(u^4)D_{40}(u^8) \\
&\quad + s^2(2L_{40}(u^8) - L_{40}(u^1))D_{40}(u^4) \\
YM_{18} &= s(-\mathbf{i}\partial_2 + \partial_3 - \mathbf{i}\partial_4)D_{40}(u^8) \\
&\quad + 2s^2\mathbf{i}L_{40}(u^6)D_{40}(u^3) + 2s^2\mathbf{i}L_{40}(u^7)D_{40}(u^4) \\
&\quad - 2s^2\mathbf{i}L_{40}(u^3)D_{40}(u^6) - 2s^2\mathbf{i}L_{40}(u^4)D_{40}(u^7)
\end{aligned}$$

where $L_{40} = -\mathbf{i}\partial_1 + \partial_2 + \mathbf{i}\partial_3 + \partial_4$, a linear differential operator. Note that $\mathbf{i}L_{40} = \partial_1 + \mathbf{i}\partial_2 - \partial_3 + \mathbf{i}\partial_4$.

Proof Formula (1.4) can be obtained via direct computation.

From this reduced form, some solutions can be obtained easily. Especially, any solution $u^i(x)$ of

$$D_{40}(u^i) = 0$$

provide a solution of Yang-Mills equation of $SU(3)$ on $R^{4,0}$. Thus we have

Proposition. By using differential linear transformation given by (1.2) and (1.3), a lot of solutions of Yang-Mills equation of $SU(3)$ on $R^{4,0}$ can be obtained from the reduced form of Yang-Mills equation

Appendix

The explicit form of the reduced form of Yang-Mills equations YM_{ij} , $i = 2, 3, 4$, $j = 1, 2, \dots, 8$ are given in this appendix.

$i = 2$ the reduced form of Yang-Mills equation is :

$$\begin{aligned}
YM_{21} &= s(\mathbf{i}\partial_1 + \mathbf{i}\partial_3 + \partial_4)D_{40}(u^1) \\
&\quad + 2s^2L_{40}(u^5)D_{40}(u^2) + 2s^2L_{40}(u^6)D_{40}(u^3) \\
&\quad - 2s^2L_{40}(u^2)D_{40}(u^5) - 2s^2L_{40}(u^3)D_{40}(u^6) \\
YM_{22} &= s(\mathbf{i}\partial_1 + \mathbf{i}\partial_3 + \partial_4)D_{40}(u^2) \\
&\quad - 2\mathbf{i}s^2L_{40}(u^5)D_{40}(u^1) + \mathbf{i}s^2L_{40}(u^7)D_{40}(u^3) \\
&\quad + \mathbf{i}s^2L_{40}(u^6)D_{40}(u^4) - \mathbf{i}s^2L_{40}(u^4)D_{40}(u^6) \\
&\quad - \mathbf{i}s^2L_{40}(u^3)D_{40}(u^7) + \mathbf{i}s^2L_{40}(u^5)D_{40}(u^8) \\
&\quad + \mathbf{i}s^2(2L_{40}(u^1) - L_{40}(u^8))D_{40}(u^5)
\end{aligned}$$

$$\begin{aligned}
YM_{23} &= s(\mathbf{i}\partial_1 + \mathbf{i}\partial_3 + \partial_4)D_{40}(u^3) \\
&\quad -\mathbf{i}s^2L_{40}(u^6)D_{40}(u^1) - \mathbf{i}s^2L_{40}(u^7)D_{40}(u^2) \\
&\quad +\mathbf{i}s^2L_{40}(u^5)D_{40}(u^4) - \mathbf{i}s^2L_{40}(u^4)D_{40}(u^5) \\
&\quad +\mathbf{i}s^2L_{40}(u^2)D_{40}(u^7) - \mathbf{i}s^2L_{40}(u^6)D_{40}(u^8) \\
&\quad +\mathbf{i}s^2(L_{40}(u^1) + L_{40}(u^8))D_{40}(u^6) \\
YM_{24} &= s(\mathbf{i}\partial_1 + \mathbf{i}\partial_3 + \partial_4)D_{40}(u^4) \\
&\quad +\mathbf{i}s^2L_{40}(u^7)D_{40}(u^1) - \mathbf{i}s^2L_{40}(u^6)D_{40}(u^2) \\
&\quad -\mathbf{i}s^2L_{40}(u^5)D_{40}(u^3) + \mathbf{i}s^2L_{40}(u^3)D_{40}(u^5) \\
&\quad +\mathbf{i}s^2L_{40}(u^2)D_{40}(u^6) - 2\mathbf{i}s^2L_{40}(u^7)D_{40}(u^8) \\
&\quad +\mathbf{i}s^2(2L_{40}(u^8) - L_{40}(u^1))D_{40}(u^7) \\
YM_{25} &= s(\mathbf{i}\partial_1 + \mathbf{i}\partial_3 + \partial_4)D_{40}(u^5) \\
&\quad +2\mathbf{i}s^2L_{40}(u^2)D_{40}(u^1) + \mathbf{i}s^2L_{40}(u^4)D_{40}(u^3) \\
&\quad -\mathbf{i}s^2L_{40}(u^3)D_{40}(u^4) + \mathbf{i}s^2L_{40}(u^7)D_{40}(u^6) \\
&\quad -\mathbf{i}s^2L_{40}(u^6)D_{40}(u^7) - \mathbf{i}s^2L_{40}(u^2)D_{40}(u^8) \\
&\quad -\mathbf{i}s^2(2L_{40}(u^1) - L_{40}(u^8))D_{40}(u^2) \\
YM_{26} &= s(\mathbf{i}\partial_1 + \mathbf{i}\partial_3 + \partial_4)D_{40}(u^6) \\
&\quad +\mathbf{i}s^2L_{40}(u^3)D_{40}(u^1) + \mathbf{i}s^2L_{40}(u^4)D_{40}(u^2) \\
&\quad -\mathbf{i}s^2L_{40}(u^2)D_{40}(u^4) - \mathbf{i}s^2L_{40}(u^7)D_{40}(u^5) \\
&\quad +\mathbf{i}s^2L_{40}(u^5)D_{40}(u^7) + \mathbf{i}s^2L_{40}(u^3)D_{40}(u^8) \\
&\quad -\mathbf{i}s^2(L_{40}(u^1) + L_{40}(u^8))D_{40}(u^3) \\
YM_{27} &= s(\mathbf{i}\partial_1 + \mathbf{i}\partial_3 + \partial_4)D_{40}(u^7) \\
&\quad -\mathbf{i}s^2L_{40}(u^4)D_{40}(u^1) + \mathbf{i}s^2L_{40}(u^3)D_{40}(u^2) \\
&\quad -\mathbf{i}s^2L_{40}(u^2)D_{40}(u^3) + \mathbf{i}s^2L_{40}(u^6)D_{40}(u^5) \\
&\quad -\mathbf{i}s^2L_{40}(u^5)D_{40}(u^6) + 2\mathbf{i}s^2L_{40}(u^4)D_{40}(u^8) \\
&\quad -\mathbf{i}s^2(2L_{40}(u^8) - L_{40}(u^1))D_{40}(u^4) \\
YM_{28} &= s(\mathbf{i}\partial_1 + \mathbf{i}\partial_3 + \partial_4)D_{40}(u^8) \\
&\quad +2s^2L_{40}(u^6)D_{40}(u^3) + 2s^2L_{40}(u^7)D_{40}(u^4) \\
&\quad -2s^2L_{40}(u^3)D_{40}(u^6) - 2s^2L_{40}(u^4)D_{40}(u^7)
\end{aligned}$$

where $L_{40} = -\mathbf{i}\partial_1 + \partial_2 - \mathbf{i}\partial_3 - \partial_4$.

$i = 3$ the reduced form of Yang-Mills equation is :

$$\begin{aligned}
YM_{31} &= s(\partial_1 + \mathbf{i}\partial_2 - \mathbf{i}\partial_4)D_{40}(u^1) \\
&\quad -2\mathbf{i}s^2L_{40}(u^5)D_{40}(u^2) - 2\mathbf{i}s^2L_{40}(u^6)D_{40}(u^3) \\
&\quad +2\mathbf{i}s^2L_{40}(u^2)D_{40}(u^5) + 2\mathbf{i}s^2L_{40}(u^3)D_{40}(u^6) \\
YM_{32} &= s(\partial_1 + \mathbf{i}\partial_2 - \mathbf{i}\partial_4)D_{40}(u^2) \\
&\quad -2s^2L_{40}(u^5)D_{40}(u^1) + s^2L_{40}(u^7)D_{40}(u^3) \\
&\quad +s^2L_{40}(u^6)D_{40}(u^4) - s^2L_{40}(u^4)D_{40}(u^6) \\
&\quad -s^2L_{40}(u^3)D_{40}(u^7) + s^2L_{40}(u^5)D_{40}(u^8) \\
&\quad +s^2(2L_{40}(u^1) - L_{40}(u^8))D_{40}(u^5)
\end{aligned}$$

$$\begin{aligned}
YM_{33} &= s(\partial_1 + \mathbf{i}\partial_2 - \mathbf{i}\partial_4)D_{40}(u^3) \\
&\quad - s^2L_{40}(u^6)D_{40}(u^1) - s^2L_{40}(u^7)D_{40}(u^2) \\
&\quad + s^2L_{40}(u^5)D_{40}(u^4) - s^2L_{40}(u^4)D_{40}(u^5) \\
&\quad + s^2L_{40}(u^2)D_{40}(u^7) - s^2L_{40}(u^6)D_{40}(u^8) \\
&\quad + s^2(L_{40}(u^1) + L_{40}(u^8))D_{40}(u^6) \\
YM_{34} &= s(\partial_1 + \mathbf{i}\partial_2 - \mathbf{i}\partial_4)D_{40}(u^4) \\
&\quad + \mathbf{i}s^2L_{40}(u^7)D_{40}(u^1) - \mathbf{i}s^2L_{40}(u^6)D_{40}(u^2) \\
&\quad - \mathbf{i}s^2L_{40}(u^5)D_{40}(u^3) + \mathbf{i}s^2L_{40}(u^3)D_{40}(u^5) \\
&\quad + \mathbf{i}s^2L_{40}(u^2)D_{40}(u^6) - 2\mathbf{i}s^2L_{40}(u^7)D_{40}(u^8) \\
&\quad + \mathbf{i}s^2(2L_{40}(u^8) - L_{40}(u^1))D_{40}(u^7) \\
YM_{35} &= s(\partial_1 + \mathbf{i}\partial_3 - \mathbf{i}\partial_4)D_{40}(u^5) \\
&\quad + 2s^2L_{40}(u^2)D_{40}(u^1) + s^2L_{40}(u^4)D_{40}(u^3) \\
&\quad - s^2L_{40}(u^3)D_{40}(u^4) + s^2L_{40}(u^7)D_{40}(u^6) \\
&\quad - s^2L_{40}(u^6)D_{40}(u^7) - s^2L_{40}(u^2)D_{40}(u^8) \\
&\quad - s^2(2L_{40}(u^1) - L_{40}(u^8))D_{40}(u^2) \\
YM_{36} &= s(\partial_1 + \mathbf{i}\partial_2 + \mathbf{i}\partial_4)D_{40}(u^6) \\
&\quad + s^2L_{40}(u^3)D_{40}(u^1) + s^2L_{40}(u^4)D_{40}(u^2) \\
&\quad - s^2L_{40}(u^2)D_{40}(u^4) - s^2L_{40}(u^7)D_{40}(u^5) \\
&\quad + s^2L_{40}(u^5)D_{40}(u^7) + s^2L_{40}(u^3)D_{40}(u^8) \\
&\quad - s^2(L_{40}(u^1) + L_{40}(u^8))D_{40}(u^3) \\
YM_{37} &= s(\partial_1 + \mathbf{i}\partial_3 - \mathbf{i}\partial_4)D_{40}(u^7) \\
&\quad - s^2L_{40}(u^4)D_{40}(u^1) + s^2L_{40}(u^3)D_{40}(u^2) \\
&\quad - s^2L_{40}(u^2)D_{40}(u^3) + s^2L_{40}(u^6)D_{40}(u^5) \\
&\quad - s^2L_{40}(u^5)D_{40}(u^6) + 2s^2L_{40}(u^4)D_{40}(u^8) \\
&\quad - s^2(2L_{40}(u^8) - L_{40}(u^1))D_{40}(u^4) \\
YM_{38} &= s(\partial_1 + \mathbf{i}\partial_3 - \mathbf{i}\partial_4)D_{40}(u^8) \\
&\quad - 2\mathbf{i}s^2L_{40}(u^6)D_{40}(u^3) - 2\mathbf{i}s^2L_{40}(u^7)D_{40}(u^4) \\
&\quad + 2\mathbf{i}s^2L_{40}(u^3)D_{40}(u^6) + 2\mathbf{i}s^2L_{40}(u^4)D_{40}(u^7)
\end{aligned}$$

where $L_{40} = -\mathbf{i}\partial_1 + \partial_2 - \mathbf{i}\partial_3 - \partial_4$.

$i = 4$ the reduced form of Yang-Mills equation is :

$$\begin{aligned}
YM_{41} &= s(\mathbf{i}\partial_1 - \partial_2 - \mathbf{i}\partial_3)D_{40}(u^1) \\
&\quad + 2s^2K_{40}(u^5)D_{40}(u^2) + 2s^2K_{40}(u^6)D_{40}(u^3) \\
&\quad - 2s^2K_{40}(u^2)D_{40}(u^5) - 2s^2K_{40}(u^3)D_{40}(u^6) \\
YM_{42} &= s(\mathbf{i}\partial_1 - \partial_2 - \mathbf{i}\partial_3)D_{40}(u^2) \\
&\quad - 2\mathbf{i}s^2K_{40}(u^5)D_{40}(u^1) + \mathbf{i}s^2K_{40}(u^7)D_{40}(u^3) \\
&\quad + \mathbf{i}s^2K_{40}(u^6)D_{40}(u^4) - \mathbf{i}s^2K_{40}(u^4)D_{40}(u^6) \\
&\quad - \mathbf{i}s^2K_{40}(u^3)D_{40}(u^7) + \mathbf{i}s^2K_{40}(u^5)D_{40}(u^8) \\
&\quad + \mathbf{i}s^2(2K_{40}(u^1) - K_{40}(u^8))D_{40}(u^5)
\end{aligned}$$

$$\begin{aligned}
YM_{43} &= s(\mathbf{i}\partial_1 - \partial_2 - \mathbf{i}\partial_3)D_{40}(u^3) \\
&\quad - \mathbf{i}s^2K_{40}(u^6)D_{40}(u^1) - \mathbf{i}s^2K_{40}(u^7)D_{40}(u^2) \\
&\quad + \mathbf{i}s^2K_{40}(u^5)D_{40}(u^4) - \mathbf{i}s^2K_{40}(u^4)D_{40}(u^5) \\
&\quad + \mathbf{i}s^2K_{40}(u^2)D_{40}(u^7) - \mathbf{i}s^2K_{40}(u^6)D_{40}(u^8) \\
&\quad + \mathbf{i}s^2(K_{40}(u^1) + K_{40}(u^8))D_{40}(u^6) \\
YM_{44} &= s(\mathbf{i}\partial_1 - \partial_2 - \mathbf{i}\partial_3)D_{40}(u^4) \\
&\quad + \mathbf{i}s^2K_{40}(u^7)D_{40}(u^1) - \mathbf{i}s^2K_{40}(u^6)D_{40}(u^2) \\
&\quad - \mathbf{i}s^2K_{40}(u^5)D_{40}(u^3) + \mathbf{i}s^2K_{40}(u^3)D_{40}(u^5) \\
&\quad + \mathbf{i}s^2K_{40}(u^2)D_{40}(u^6) - 2\mathbf{i}s^2K_{40}(u^7)D_{40}(u^8) \\
&\quad + \mathbf{i}s^2(2K_{40}(u^8) - K_{40}(u^1))D_{40}(u^7) \\
YM_{45} &= s(\mathbf{i}\partial_1 - \partial_2 - \mathbf{i}\partial_3)D_{40}(u^5) \\
&\quad + 2\mathbf{i}s^2K_{40}(u^2)D_{40}(u^1) + \mathbf{i}s^2K_{40}(u^4)D_{40}(u^3) \\
&\quad - \mathbf{i}s^2K_{40}(u^3)D_{40}(u^4) + \mathbf{i}s^2K_{40}(u^7)D_{40}(u^6) \\
&\quad - \mathbf{i}s^2K_{40}(u^6)D_{40}(u^7) - \mathbf{i}s^2K_{40}(u^2)D_{40}(u^8) \\
&\quad - \mathbf{i}s^2(2K_{40}(u^1) - K_{40}(u^8))D_{40}(u^2) \\
YM_{46} &= s(\mathbf{i}\partial_1 - \partial_2 - \mathbf{i}\partial_3)D_{40}(u^6) \\
&\quad + \mathbf{i}s^2K_{40}(u^3)D_{40}(u^1) + \mathbf{i}s^2K_{40}(u^4)D_{40}(u^2) \\
&\quad - \mathbf{i}s^2K_{40}(u^2)D_{40}(u^4) - \mathbf{i}s^2K_{40}(u^7)D_{40}(u^5) \\
&\quad + \mathbf{i}s^2K_{40}(u^5)D_{40}(u^7) + \mathbf{i}s^2K_{40}(u^3)D_{40}(u^8) \\
&\quad - \mathbf{i}s^2(K_{40}(u^1) + K_{40}(u^8))D_{40}(u^3) \\
YM_{47} &= s(\mathbf{i}\partial_1 - \partial_2 - \mathbf{i}\partial_3)D_{40}(u^7) \\
&\quad - \mathbf{i}s^2K_{40}(u^4)D_{40}(u^1) + \mathbf{i}s^2K_{40}(u^3)D_{40}(u^2) \\
&\quad - \mathbf{i}s^2K_{40}(u^2)D_{40}(u^3) + \mathbf{i}s^2K_{40}(u^6)D_{40}(u^5) \\
&\quad - \mathbf{i}s^2K_{40}(u^5)D_{40}(u^6) + 2\mathbf{i}s^2K_{40}(u^4)D_{40}(u^8) \\
&\quad - \mathbf{i}s^2(2K_{40}(u^8) - K_{40}(u^1))D_{40}(u^4) \\
YM_{48} &= s(\mathbf{i}\partial_1 - \partial_2 - \mathbf{i}\partial_3)D_{40}(u^8) \\
&\quad + 2s^2K_{40}(u^6)D_{40}(u^3) + 2s^2K_{40}(u^7)D_{40}(u^4) \\
&\quad - 2s^2K_{40}(u^3)D_{40}(u^6) - 2s^2K_{40}(u^4)D_{40}(u^7)
\end{aligned}$$

where $K_{40} = -\mathbf{i}\partial_1 + \partial_2 + \mathbf{i}\partial_3 + \partial_4$.

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